

Cold atoms 3: Optical lattices

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States that are *not* simple Bose-Einstein condensates:

- Bose glass
- Mott insulators

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Optical potentials

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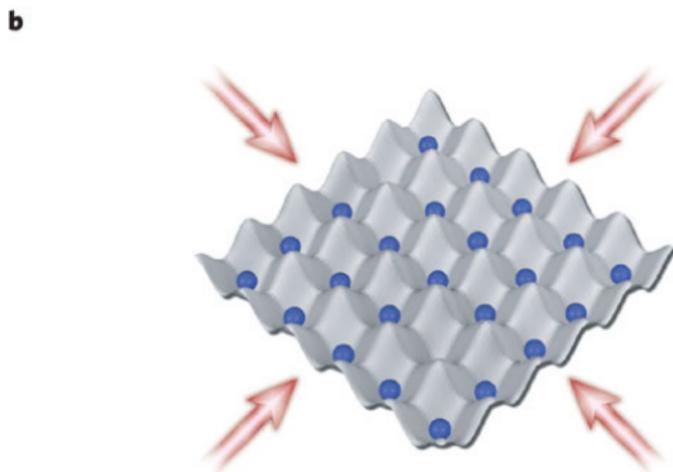
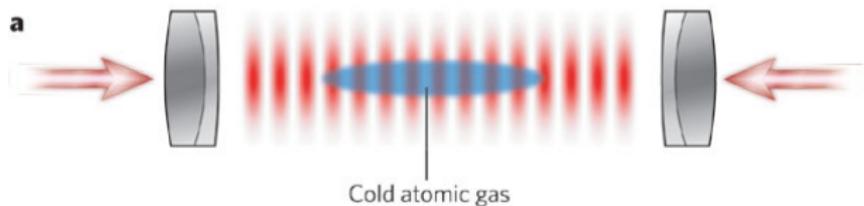
Nevertheless, an E-field induces a *dipole moment*.

Second-order perturbation theory:

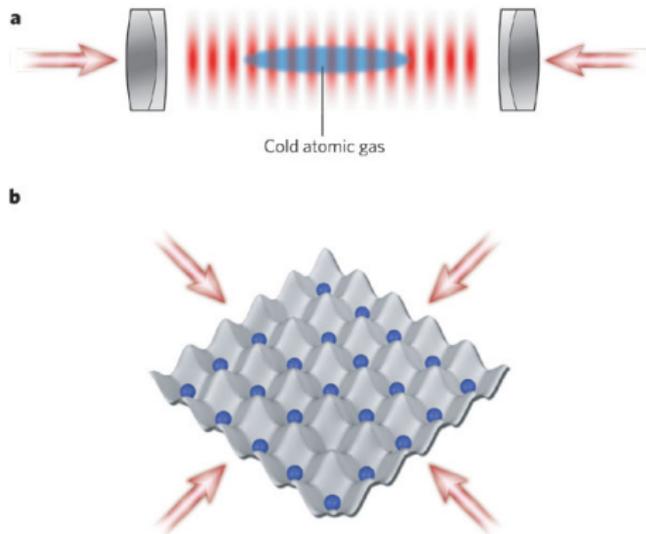
$$V(r) \propto |E(r)|^2 \propto I(r)$$

I is the irradiance

Standing wave



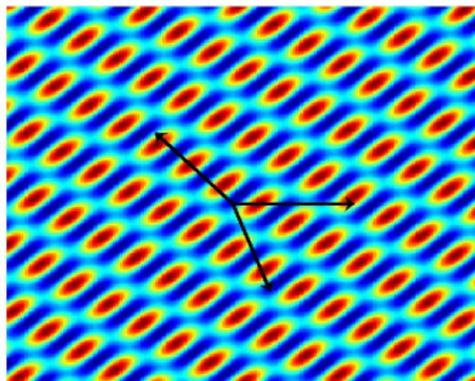
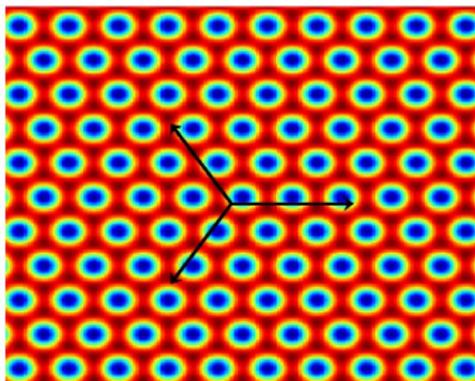
Standing wave



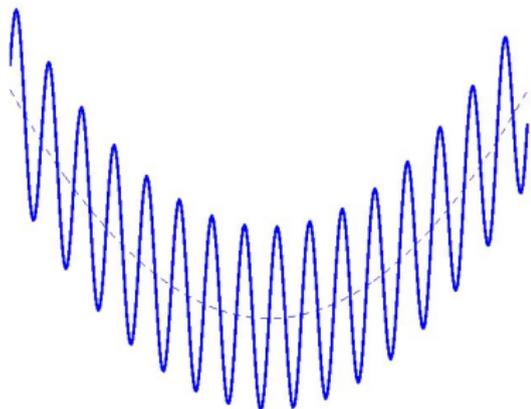
Optical lattice: $V(r) = V_{01} \cos \vec{k}_1 \cdot \vec{r} + V_{02} \cos \vec{k}_2 \cdot \vec{r}$

Optical lattices

Generally: Several laser beams create *interference patterns* in 1, 2, or 3 dimensions.



Lattice plus trap



Optical lattice: periodicity
Magnetic trap: Confinement

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- Fundamental interest: Quantum mechanics + random potential
- New quantum phase: The Bose glass

Disorder with optical potentials

Alternative 1: Laser speckle
Light diffracted through roughened glass



gives a random pattern

Alternative 2: Quasiperiodic potential

Disorder with optical potentials

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Simple 1D pattern:

$$V(x) = A \cos k_1 x + B \cos k_2 x$$

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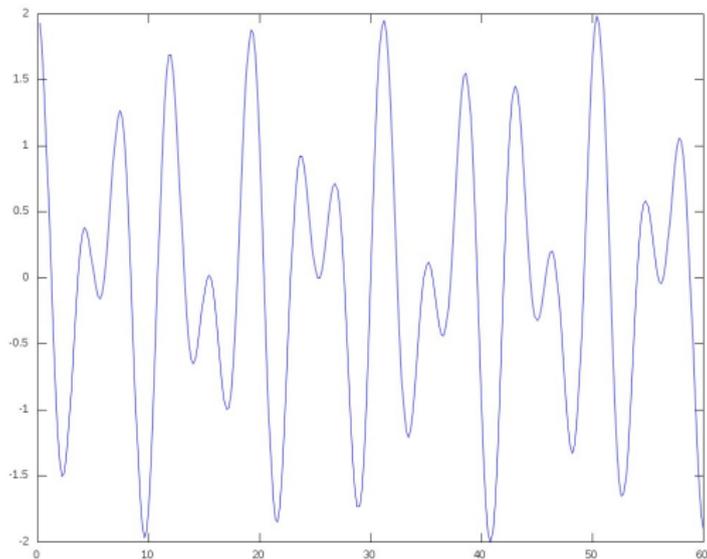
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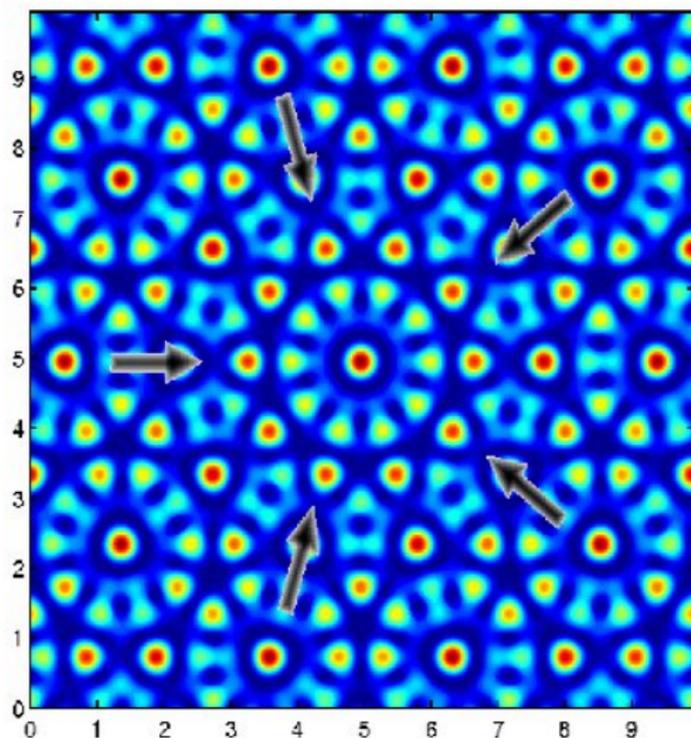
(In practice: make sure that k_1/k_2 is not close to a ratio of small integers such as $2/3$ or $1/4$.)

Popular choice: The golden ratio $k_1/k_2 = (1 + \sqrt{5})/2$

Quasiperiodic potential in 1D



Quasiperiodic potential in 2D



Anderson localization

A quantum particle in a random potential will not propagate:
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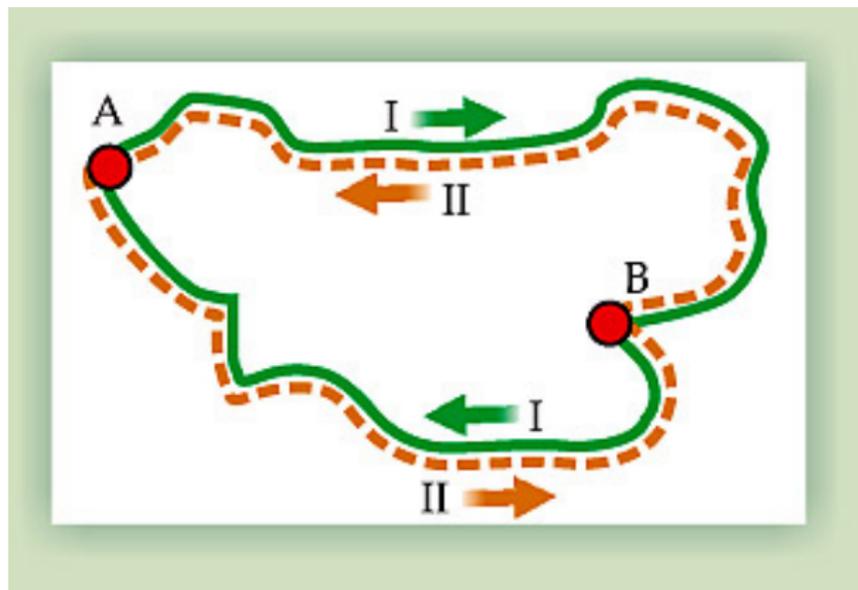
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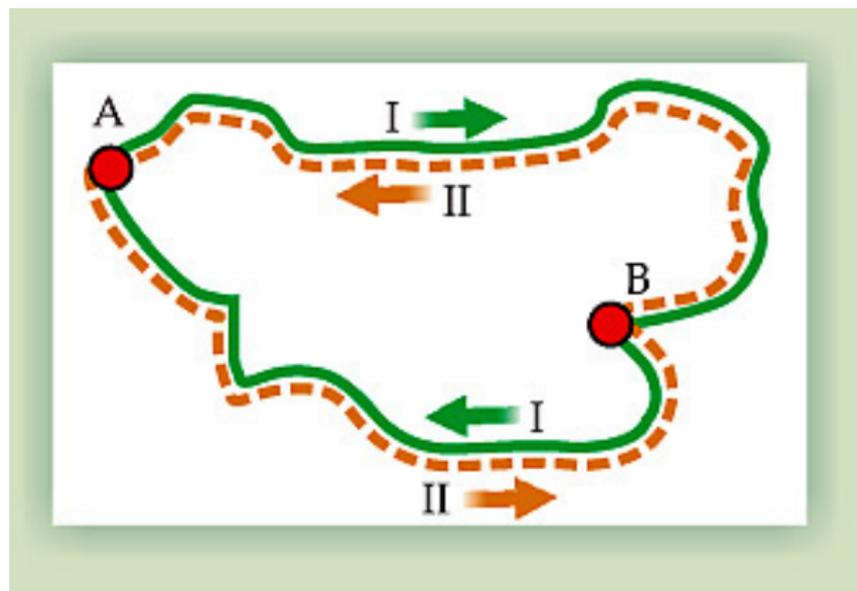
Non-periodic potentials: Interference effects \rightarrow localization.

Anderson's argument



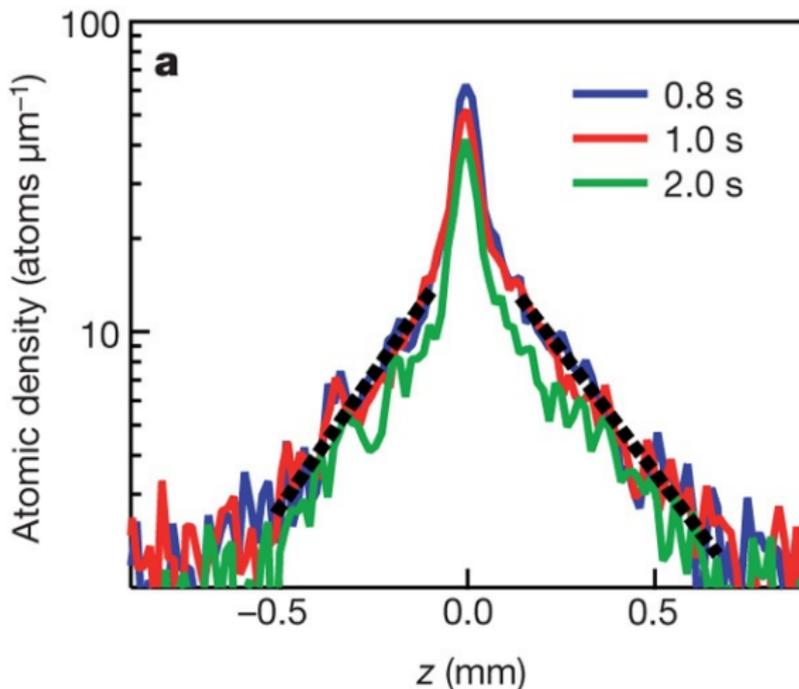
Constructive interference between a path and its reverse.

Anderson's argument



Constructive interference between a path and its reverse.
If not periodic, then *no* constructive interference for *other* paths –
only those that lead back to point of origin!

Exponential localization

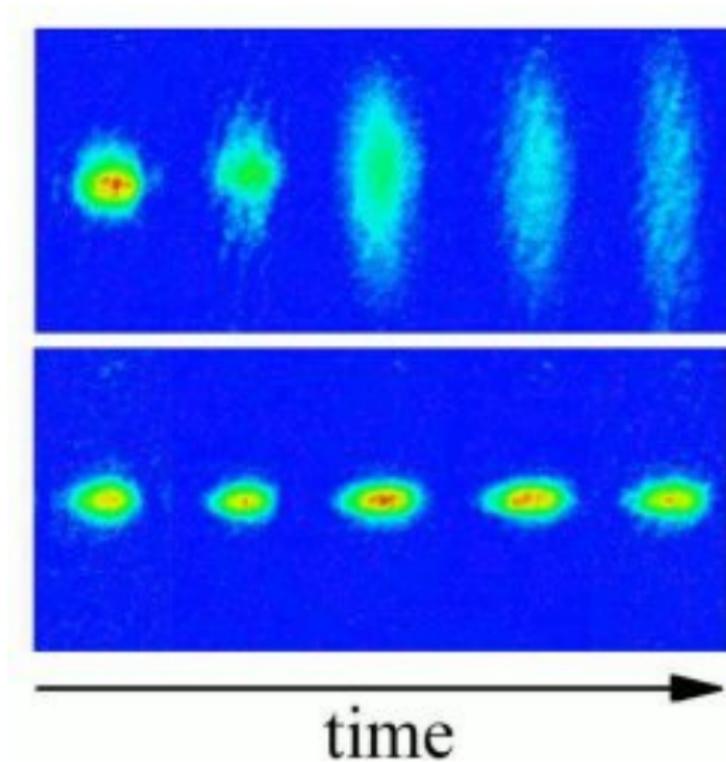


Wavefunctions decay exponentially at long distances:
Localized eigenstates.

Observation of Anderson localization

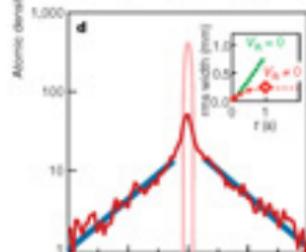
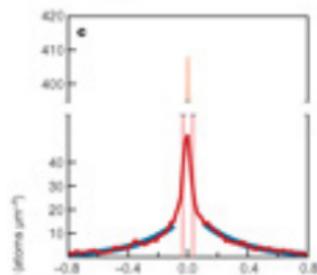
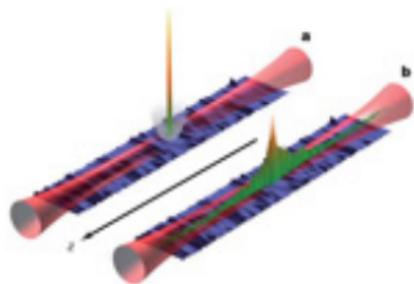
- Not in solids: Mean free path too short
- Must construct artificial system in order to see A.L.
- Light: 1997
- BEC: 2008
- Very dilute BEC: neglect interactions.
- AL is a single-particle phenomenon.

Experiment, Florence



Roati *et al.*, Florence. Quasiperiodic potential in 1D

Experiment, Paris



Billy *et al.*, Paris.
Speckle potential in 1D

What about interacting Bose gases?

- A quantum phase was predicted, the “Bose glass” phase.

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- For the experts: Gapless but not superfluid.
- Loosely speaking: System is broken up into several small BEC, whose phases are mutually random

Define the particle-particle correlation function

$$g(r, r') = \int dr_2 \cdots dr_N \psi^*(r, r_2, \dots, r_N) \psi(r', r_2, \dots, r_N)$$

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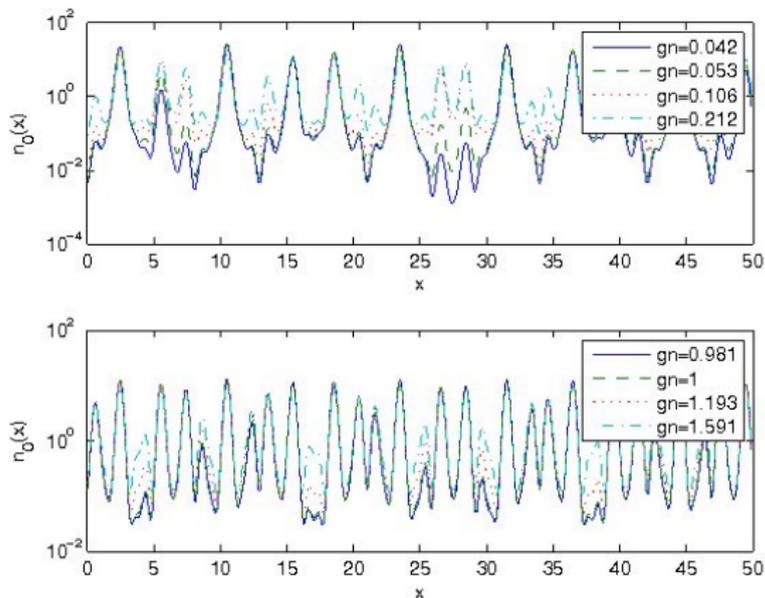
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BEC: $g(r, r') \sim \text{const}$ at large $r - r'$

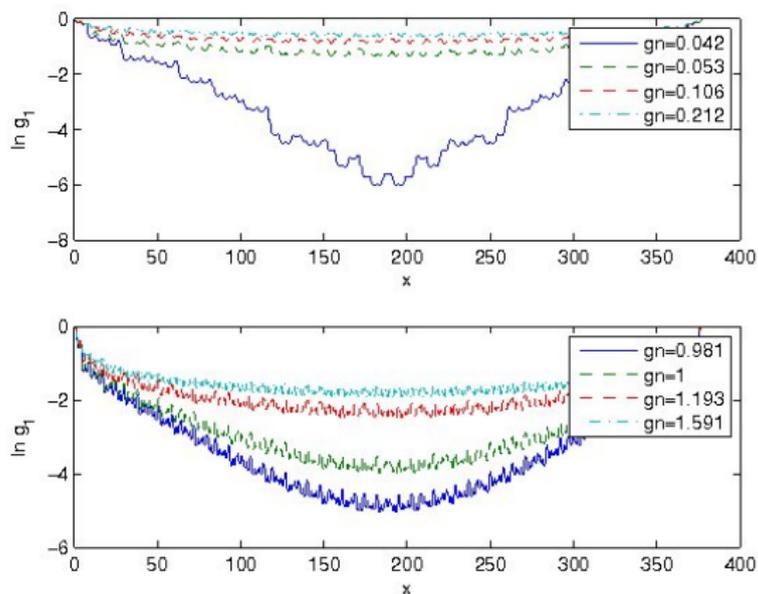
Bose glass: $g(r, r') \sim e^{-|r-r'|}$ exponential decay

Bose glass: Density



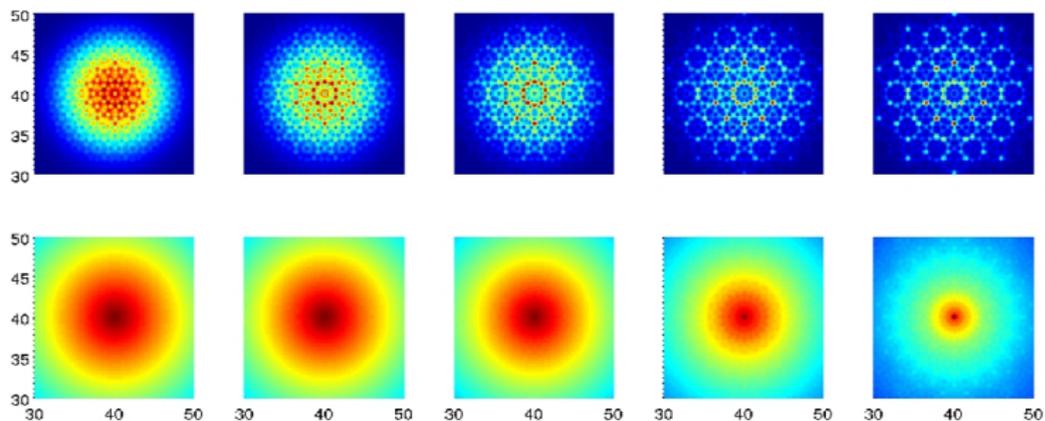
Cetoli and Lundh, 2010

Bose glass: Correlations



Cetoli and Lundh, 2010

Bose glass in 2D



Cetoli and Lundh, to be published

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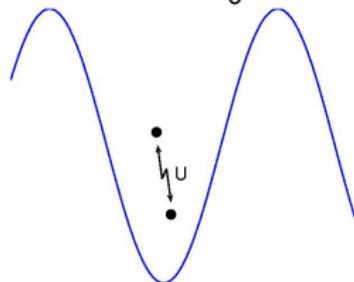
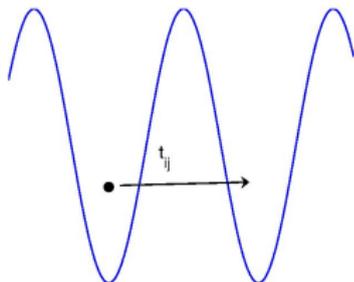
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Put all this together into Hamiltonian:

$$H - \mu N = -t \sum_{\langle i,j \rangle} a_i^\dagger a_j + \frac{U}{2} \sum_i a_i^\dagger a_i^\dagger a_i a_i - \mu \sum_i a_i^\dagger a_i$$

Mott transition

Bosonic Hubbard model has two phases:

- Superfluid: When kinetic energy dominates, J/U not small.
A BEC forms.

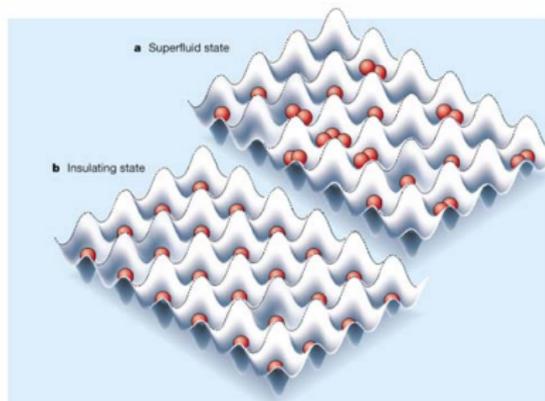
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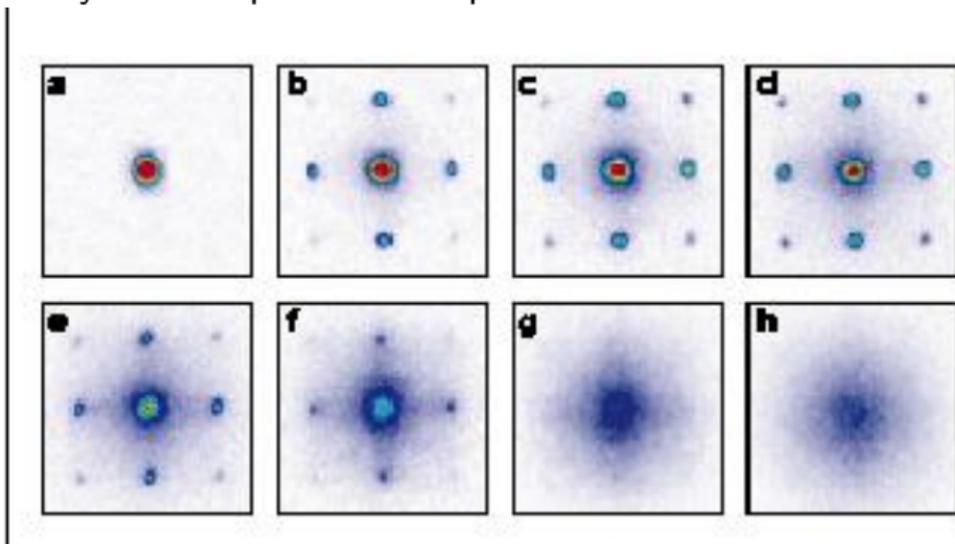
Mott transition

- A quantum phase transition.
- Takes place at zero temperature.
- Driven by quantum fluctuations, not thermal.
- Can be described by:
 - Monte Carlo
 - DMRG
 - Mean-field theories
 - Strong-coupling expansion
 - some more
- Mostly heavy numerical computations.

Observing the Mott transition

Key observation: BEC is coherent – will give interference fringes.
Mott insulator is not.

Periodic system \Rightarrow peaks at reciprocal lattice vectors!



First experiment by Greiner et al., 2001

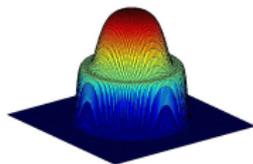
Trapped systems

$$H - \mu N = -t \sum_{\langle i,j \rangle} a_i^\dagger a_j + \frac{U}{2} \sum_i a_i^\dagger a_i^\dagger a_i a_i + \sum_i \left(\frac{1}{2} \omega^2 r_i^2 - \mu \right) a_i^\dagger a_i$$

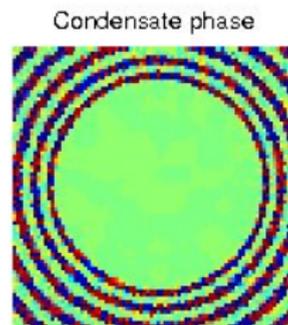
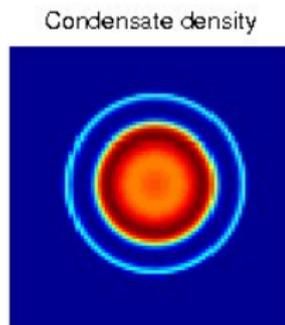
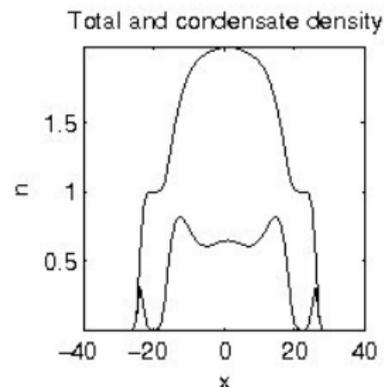
Superfluid and Mott insulating regions coexist

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Mott and superfluid shells in trap



Density – condensate density – condensate phase
Note that a *local-density approximation* applies.

Rotation in optical lattice

Can we have vortices in an optical lattice?

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Techniques:

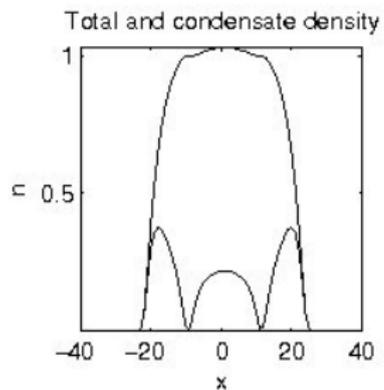
- Rotate the actual lattice
 - Mechanically rotating mask
 - Interference pattern modulated in time
- Artificial gauge fields

What I found

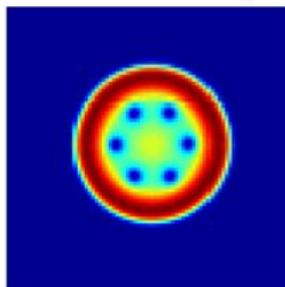
Two limits:

1. Close to the Mott transition; vortex has Mott insulating atoms in the core
2. Further from the Mott transition; vortex core is smaller than the lattice spacing. Vortex sits between the lattice sites where density vanishes

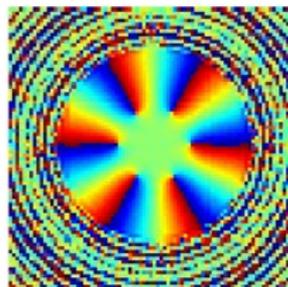
Vortices with Mott cores



Condensate density

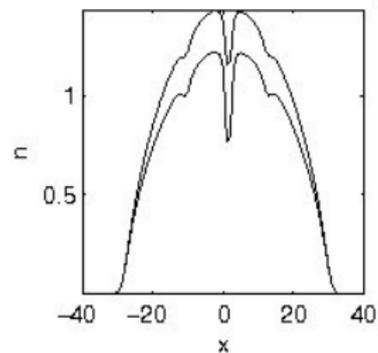


Condensate phase

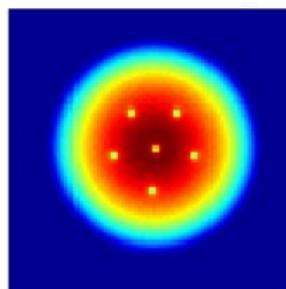


Interstitial vortices

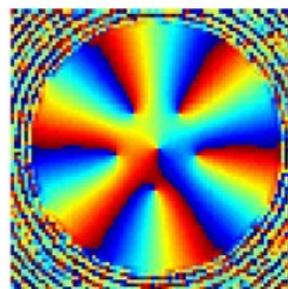
Total and condensate density



Condensate density



Condensate phase



End of lecture 3

Thank you for your attention